

Motions in a Riemannian space and some integrability conditions

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Abstract

Knebelman¹ studied and defined collineations and motions in generalised spaces. Levine³ studied motions in linearly connected two dimensional spaces. Takano², studied on the existence of affine motion in a space with recurrent curvature tensor. Further, Negi and Rawat⁶ studied affine motion in an almost Tachibana recurrent space. Rawat and Dobhal⁹, studied on Projective motion in a Tachibana symmetric space

In the present paper, we have studied Affine motion, Projective motion, Conformal motion, Integrability conditions of Killing equations also several theorems have been established and proved therein.

1. Motions in a Riemannian space.

Consider an n-dimensional Riemannian space V_n covered by a set of neighbourhoods with Coordinates ξ^x and endowed with the fundamental quadratic differential form

$$ds^2 = g_{\lambda x}(\xi) d\xi^\lambda d\xi^x \quad (1.1)$$

where the Greek indices $x, \lambda, \mu, \nu, \dots \dots$, run over the range 1, 2, 3, ... n.

In the V_n referred to ξ^x , we consider a point transformation

$$T : \quad * \xi^x = f^x(\xi^v); \text{Det}(\partial_\lambda f^x) \neq 0 \quad (1.2)$$

which establishes a one-to-one correspondence between the points of a region R and those of some other region $*R$, where ∂_λ stands for the partial derivation $\frac{\partial}{\partial \xi^\lambda}$.

During this point transformation, a point ξ^x in R is carried to a point $*\xi^x$ in $*R$ and a point $\xi^x + d\xi^x$ in R to a point $*\xi^x + d*\xi^x$ in $*R$.

If the distance $d*s$ between two displaced points $*\xi^x$ and $*\xi^x + d*\xi^x$ is always

equal to the distance between the two original points ξ^x and $\xi^x + d\xi^x$ the point transformation (1.2) is called a motion or an isometry in the V_n .

(i) *Affine motion in V_n* : When a point transformation (1.2), transforms any pair of parallel vectors into a pair of parallel vectors, then (1.2) is called Affine motion in a V_n .

For an affine motion, we must have

$$\delta^m u^x (*\xi) \stackrel{\text{def}}{=} d^m u^x (*\xi) + \Gamma_{\mu\lambda}^x (*\xi) \frac{m^\lambda}{u} (*\xi) d*\xi^\mu = 0 \quad (1.3)$$

(ii) *Projective Motions in V_n* : When a point transformation (1.2) transforms the system of geodesics into the same system, then (1.2) is called a projective Motion in V_n

The necessary and sufficient condition that (1.2) be a projective motion in a V_n is that the Lie-difference of $\Gamma_{\mu\lambda}^x$ with respect to (1.2) has the form

$$*\Gamma_{\mu\lambda}^x - \Gamma_{\mu\lambda}^x = A_\mu^x p_\lambda + A_\lambda^x p_\mu, \quad (1.4)$$

where p_λ is a covariant vector.

When (1.2) is an infinitesimal transformation $*\xi^x = \xi^x + v^x(\xi)dt$, then the condition is

$$\frac{\mathfrak{L}}{v} \Gamma_{\mu\lambda}^x = A_\mu^x p_\lambda + A_\lambda^x p_\mu \quad (1.5)$$

(iii) *Conformal Motion in V_n* : When a point transformation (1.2) does not change the angle between two direction at a point, then (1.2) is called a conformal motion in V_n . The necessary and sufficient condition that (1.2) be conformal motion in a V_n is that the Lie -difference of $\mathcal{G}_{\lambda x}$ with respect to (1.2) be proportional to $\mathcal{G}_{\lambda x}$. [Schouten]

$$'\mathcal{G}_{\lambda x} - \mathcal{G}_{\lambda x} = 2\phi \mathcal{G}_{\lambda x} \quad (1.6)$$

Where ϕ is a scalar

When (1.2) is an infinitesimal transformation $*\xi^x = f^x + v^x(\xi)dt$, then the condition is

$$\frac{\mathfrak{L}}{v} \mathcal{G}_{\lambda x} = 2\phi \mathcal{G}_{\lambda x}. \quad (1.7)$$

Now, we have the following theorem which is geometrically evident.

Theorem (1.1) : A motion in a V_n is an Affine motion.

Proof : To prove this, we apply the formula

$$\frac{\mathfrak{L}}{v} \nabla_\nu p_\mu^{x\lambda} - \nabla_\nu \frac{\mathfrak{L}}{v} p_{\dots\mu}^{x\lambda} = \left(\frac{\mathfrak{L}}{v} \Gamma_{\nu\rho}^x \right) P_{\dots\mu}^{\rho\lambda} + \left(\frac{\mathfrak{L}}{v} \Gamma_{\nu\rho}^\lambda \right) P_{\dots\mu}^{x\rho} - \left(\frac{\mathfrak{L}}{v} \Gamma_{\nu\mu}^\rho \right) P_{\dots\rho}^{x\lambda} \quad (1.8)$$

to the fundamental tensor $\mathcal{G}_{\lambda x}$, we have

$$\begin{aligned} & \frac{\mathfrak{L}}{v} (\nabla_\mu \mathcal{G}_{\lambda x}) - \nabla_\mu \left(\frac{\mathfrak{L}}{v} \mathcal{G}_{\lambda x} \right) \\ &= - \left(\frac{\mathfrak{L}}{v} \left\{ \begin{matrix} \rho \\ \mu \lambda \end{matrix} \right\} \right) \mathcal{G}_{\rho x} - \left(\frac{\mathfrak{L}}{v} \left\{ \begin{matrix} \rho \\ \mu x \end{matrix} \right\} \right) \mathcal{G}_{\lambda\rho}, \end{aligned} \quad (1.9)$$

from which

$$\frac{\mathfrak{L}}{v} \left\{ \begin{matrix} x \\ \mu \lambda \end{matrix} \right\} = \frac{1}{2} g^{x\rho} \left[\nabla_\mu \frac{\mathfrak{L}}{v} \mathcal{G}_{\lambda\rho} + \nabla_\lambda \frac{\mathfrak{L}}{v} \mathcal{G}_{\mu\rho} - \nabla_\rho \frac{\mathfrak{L}}{v} \mathcal{G}_{\mu\lambda} \right]. \quad (1.10)$$

This equation shows that $\frac{\mathfrak{L}}{v} \mathcal{G}_{\lambda x} = 0$

implies $\frac{\mathfrak{L}}{v} \left\{ \begin{matrix} x \\ \mu \lambda \end{matrix} \right\} = 0$,

[**Note** : Under some global conditions

$\frac{\mathfrak{L}}{v} \left\{ \begin{matrix} x \\ \mu \lambda \end{matrix} \right\} = 0$ implies $\frac{\mathfrak{L}}{v} \mathcal{G}_{\lambda x} = 0$.]

Theorem (1.2) : For a motion in a V_n the Lie-derivative of the curvature tensor and

its successive covariant derivative Vanish.

Proof: To prove the above theorem, we apply the following

$$\begin{aligned} \nabla_\nu \mathfrak{L} \Gamma_{\mu\lambda}^x - \nabla_\mu \mathfrak{L} \Gamma_{\nu\lambda}^x + 2 S_{\nu\mu}^{\dots\rho} \mathfrak{L} \Gamma_{\rho\lambda}^x \\ = \mathfrak{L} R_{\nu\mu\lambda}^x \end{aligned} \quad (1.11)$$

to the Christoffel symbol, we have

$$\nabla_\nu \mathfrak{L} \left\{ \begin{matrix} x \\ \mu \lambda \end{matrix} \right\} - \nabla_\mu \mathfrak{L} \left\{ \begin{matrix} x \\ \nu \lambda \end{matrix} \right\} = \mathfrak{L} K_{\nu\mu\lambda}^x.$$

Where $K_{\nu\mu\lambda}^x$ is the curvature tensor of V_n .

Thus for a motion, we have

$$\mathfrak{L} K_{\nu\mu\lambda}^x = 0. \quad (1.12)$$

On the other hand since a motion is an affine motion, the covariant derivation and the Lie-derivation are Commutative. Thus from (1.12), we obtain

$$\mathfrak{L} \nabla_\omega K_{\nu\mu\lambda}^x = 0, \mathfrak{L} \nabla_{\omega_2} \nabla_{\omega_1} K_{\nu\mu\lambda}^x = 0, \dots$$

This proves the theorem.

2. Theorems on Projectively or Conformally related spaces:

Consider two Riemannian spaces V_n and V_n^* which are in geodesic correspondence. Then denoting the Christoffel symbols of them

by $\left\{ \begin{matrix} x \\ \mu \lambda \end{matrix} \right\}$ and $\left\{ \begin{matrix} x \\ \mu \lambda \end{matrix} \right\}^*$ respectively, we have

$$\left\{ \begin{matrix} x \\ \mu \lambda \end{matrix} \right\}^* = \left\{ \begin{matrix} x \\ \mu \lambda \end{matrix} \right\} + A_\mu^x p_\lambda + A_\lambda^x p_\mu.$$

But since V_n and V_n^* are both Riemannian, the vector p_λ should be a gradient. Thus putting

$$p_\lambda = \frac{1}{2} \partial_\lambda \log \phi, \text{ we have}$$

$$\begin{aligned} \left\{ \begin{matrix} x \\ \mu \lambda \end{matrix} \right\}^* &= \left\{ \begin{matrix} x \\ \mu \lambda \end{matrix} \right\} + \frac{1}{2} A_\mu^x \partial_\lambda \log \phi \\ &+ \frac{1}{2} A_\lambda^x \partial_\mu \log \phi. \end{aligned} \quad (2.1)$$

We now assume that the V_n admits a

motion with symbol $\mathfrak{L} f$. Then, we have

$$\begin{aligned} \mathfrak{L} g_{\lambda x} &= \nabla_\lambda v_x + \nabla_x v_\lambda = \partial_\lambda v_x \\ &+ \partial_x v_\lambda - 2 \left\{ \begin{matrix} \rho \\ \lambda x \end{matrix} \right\} v_\rho = 0. \end{aligned}$$

Consequently on utilizing (2.1), we have

$$\begin{aligned} \mathfrak{L} g_{\lambda x} &= \partial_\lambda v_x + \partial_x v_\lambda - 2 \left[\left\{ \begin{matrix} \rho \\ \lambda x \end{matrix} \right\}^* \right. \\ &- \frac{1}{2} A_\lambda^\rho \partial_x \log \phi - \frac{1}{2} A_x^\rho \partial_\lambda \log \phi \left. \right] v_\rho \\ &= \phi^{-1} \left[\partial_\lambda (\phi v_x) + \partial_x (\phi v_\lambda) - 2 \left\{ \begin{matrix} \rho \\ \lambda x \end{matrix} \right\}^* \phi v_\rho \right] \end{aligned}$$

Thus, denoting by $g_{\lambda x}^*$ the fundamental tensor

of V_n^* and $\mathfrak{L}^* f$ the symbol defined by ϕv_x

in V_n^* , we have

$$\mathfrak{L}^* g_{\lambda x} = \phi^{-1} \mathfrak{L} g_{\lambda x}^*$$

Thus, we have.

Theorem (2.1): If two Riemannian spaces V_n and V_n^* are in geodesic correspondence and if V_n admits a group of motions, V_n^* also admits a group of motions.

Theorem (2.2): If a V_n admits a G_r of motions such that the rank of v^x in a neighborhood is equal to $r < n$, then there exist $n - r$ V_n^s , Corresponding to $n - r$ independent

solutions of $\mathfrak{L}_v \rho^2 = 0$, Which are Conformal to the given V_n and admit the same group as a group of motions.

Proof: Consider a V_n which admits an r-parameter group G_r of motions such that the rank of v^x is in a certain neighborhood a equal to $r < n$. Then we have $\mathfrak{L}_v g_{\lambda x} = 0$.

In order that a space V_n^* Conformal to V_n admit the same group G_r as a group of motions, it is necessary and sufficient that there exist a function ρ^2 such that $\mathfrak{L}_v (\rho^2 g_{\lambda x}) = 0$ or

$$\mathfrak{L}_v \rho^2 = 0. \text{ But on the other hand}$$

$$\begin{pmatrix} \mathfrak{L}_c & \mathfrak{L}_b \\ \mathfrak{L}_c & \mathfrak{L}_b \end{pmatrix} \rho^2 = C_{cb}^a \mathfrak{L}_a \rho^2$$

and consequently $\mathfrak{L}_a \rho^2 = 0$ admits $(n - r)$ independent solutions.

Note: The a-rank of the $\mathfrak{L}_a g_{\lambda x}$ is the rank of the matrix $\mathfrak{L}_a g_{\lambda x}$ where a denotes the rows and λx denotes the columns.

Theorem (2.3): In order that a G_r in X_n such that the rank of v^x in a neighborhood a is equal to $r \leq n$, can be regarded as a group of motions in a C_n , it is necessary and sufficient that the group be a subgroup of a group of Conformal transformations.

Proof: The necessity is evident.

Conversely, if the group is a subgroup of a group of conformal transformations, it is a group of conformal motions in a C_n (i.e. C_n stands for a Conformally Euclidean space). Consequently, there exists a V_n which is conformal to C_n and is itself a C_n which admits the group as a group of motions.

3. Integrability Conditions of Killing's equation :

The Integrability Conditions of killing's equation

$$\mathfrak{L}_v g_{\lambda x} = \nabla_\lambda v_x + \nabla_x v_\lambda = 0 \tag{3.1}$$

can be deduced from it, considering first the equation

$$\mathfrak{L}_v \left\{ \begin{matrix} x \\ \mu \lambda \end{matrix} \right\} = \nabla_\mu \nabla_\lambda v^x + K_{v\mu\lambda}^{...x} v^\nu = 0. \tag{3.2}$$

and next the mixed system of partial differential equations

$$\left. \begin{matrix} v_{\lambda x} + v_{x\lambda} = 0 & (v_{\lambda x} = v_\lambda^\rho g_{\rho x}) \\ \nabla_\lambda v^x = v_\lambda^x, & \nabla_\mu v_\lambda^x = -K_{v\mu\lambda}^{...x} v^\nu, \end{matrix} \right\} \tag{3.3}$$

We know that the equations (3.1) and (3.2) or the equation (3.3) have for Integrability conditions

$$\begin{matrix} \mathfrak{L}_v K_{v\mu\lambda}^{...x} = 0, & \mathfrak{L}_v \nabla_\omega K_{v\mu\lambda}^{...x} = 0, \\ \mathfrak{L}_v \nabla_{\omega_2} \nabla_{\omega_1} K_{v\mu\lambda}^{...x} = 0, & \dots\dots\dots, \end{matrix} \tag{3.4}$$

4. Theorems on Affine and Projective motions :

We consider an A_n which admits a G_r of affine motions with the infinitesimal operators $\mathfrak{L}_a f = v^\mu \partial_\mu f$ such that the

rank of v^x_a in a neighborhood is $r \leq n$. Then, we have

$$\mathfrak{L}^x_a \Gamma^x_{\mu\lambda} = 0 \quad (4.1)$$

In order that an A_n^* , projectively related to the A_n , admits the G_r as a group of affine motions, it is necessary and sufficient that there exist a covariant vector field p_λ such that

$$\mathfrak{L}^x_a (\Gamma^x_{\mu\lambda} + \mu A^x_\lambda + \lambda A^x_\mu) = 0$$

Or

$$\mathfrak{L}^x_a \lambda = 0 .$$

But this system of partial differential equation is completely integrable, hence

Theorem (4.1): If an A_n admits a G_r of affine motions with the infinitesimal operators $\mathfrak{L}^x_a f = v^x_a \partial_\mu f$ such that the rank of v^x_a in a neighborhood is $r \leq n$, there exists always in A_n^* which is (not trivially) projectively related to A_n and which admits the same G_r as a group of affine motions.

We next consider an A_n which admits a G_r of projective motions such that the rank of v^x_a in a neighborhood is $r \leq n$, Then, we have

$$\mathfrak{L}^x_a \Gamma^x_{\mu\lambda} = \mu A^x_\lambda + \lambda A^x_\mu . \quad (4.3)$$

In order that an A_n^* , projectively related to the A_n , admit the same G_r as a group of affine motions, it is necessary and sufficient that there exist a covariant vector field p_λ such that

$$\mathfrak{L}^x_a (\Gamma^x_{\mu\lambda} + \mu A^x_\lambda + \lambda A^x_\mu) = 0 \quad (4.4)$$

Or

$$\mathfrak{L}^x_a \lambda = - \frac{\lambda}{a} . \quad (4.5)$$

On the other hand, substituting (4.3) in the identity

$$\left(\mathfrak{L}^x_c \mathfrak{L}^x_b \Gamma^x_{\mu\lambda} \right) = C^a_{cb} \mathfrak{L}^x_a \Gamma^x_{\mu\lambda} , \text{ we get}$$

$$\mathfrak{L}^x_c \lambda - \mathfrak{L}^x_b \lambda = C^a_{cb} \frac{\lambda}{a} .$$

Which shows that r covariant vectors p_λ form a complete system with respect to G_r . Thus (4.5) is completely integrable and we have

Theorem (4.2): when an A_n admits a G_r of projective motions such that the rank of v^x_a is $r \leq n$, there exists an A_n^* which is (not trivially) projectively related to A_n and which admits the same G_r as a group of affine motions.

From this, we obtain

Theorem (4.3): In order that a G_r is an X_n such that the rank of v^x_a is $r \leq n$, can be regarded as a group of affine motions in a D_n , it is necessary and sufficient that the G_r be a subgroup of the ordinary projective group.

Proof: The necessity is evident. Conversely, if the group G_r is a subgroup of the ordinary projective group, it is a group of projective motions in a D_n . consequently according to Theorem(4.2), there exist an A_n which is projectively related to D_n , and is itself

a D_n and admits G_r as a group of affine motions. Thus, theorem (4.3) is proved.

5. Integrability condition of $\mathfrak{L} \Gamma_{\mu\lambda}^x = 0$.

We consider the Integrability condition of $\mathfrak{L} \Gamma_{\mu\lambda}^x = 0$, which can be written as

$$\left. \begin{aligned} \nabla_\lambda v^x &= v_\lambda^x - 2 S_{\mu\lambda}^{\dots x} v^\mu, \\ \nabla_\lambda v_\lambda^x &= -R_{\nu\mu\lambda}^{\dots x} v^\nu, \end{aligned} \right\} \quad (5.1)$$

Now from the following equations.

$$\begin{aligned} *R_{\nu\mu\lambda}^{\dots x} &= R_{\nu\mu\lambda}^{\dots x} + \left(\nabla_\nu \mathfrak{L} \Gamma_{\mu\lambda}^x - \nabla_\mu \mathfrak{L} \Gamma_{\nu\lambda}^x \right. \\ &\quad \left. + 2S_{\nu\mu}^{\dots\rho} \mathfrak{L} \Gamma_{\rho\lambda}^x \right) dt. \end{aligned}$$

and

$$\nabla_\nu \mathfrak{L} \Gamma_{\mu\lambda}^x - \nabla_\mu \mathfrak{L} \Gamma_{\nu\lambda}^x + 2S_{\nu\mu}^{\dots\rho} \mathfrak{L} \Gamma_{\rho\lambda}^x = \mathfrak{L} R_{\nu\mu\lambda}^{\dots x},$$

we have

$$\mathfrak{L} S_{\mu\lambda}^{\dots x} = 0, \quad \mathfrak{L} R_{\nu\mu\lambda}^{\dots x} = 0, \quad (5.2)$$

respectively. Then applying the formula

$$\begin{aligned} \mathfrak{L} \nabla_\nu P_{\dots\mu}^{x\lambda} - \nabla_\nu \mathfrak{L} P_{\dots\mu}^{x\lambda} &= \left(\mathfrak{L} \Gamma_{\nu\lambda}^x \right) \\ P_{\dots\mu}^{\rho\lambda} + \left(\mathfrak{L} \Gamma_{\nu\rho}^\lambda \right) P_{\dots\mu}^{x\rho} - \left(\mathfrak{L} \Gamma_{\nu\mu}^\rho \right) P_{\dots\rho}^{x\lambda}. \end{aligned}$$

to $S_{\mu\lambda}^{\dots x}$ and $R_{\nu\mu\lambda}^{\dots x}$, we obtain

$$\mathfrak{L} \nabla_\nu P_{\dots\mu}^{x\lambda} = 0, \quad \mathfrak{L} \nabla_\omega R_{\nu\mu\lambda}^{\dots x} = 0. \quad (5.3)$$

respectively.

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